

TITLE PHASE DETONATED SHOCK TUBE (PDST)

AUTHOR(S) WILLIAM D. ZERWEKH, M-6
STANLEY P. MARSH, M-6
TAI-HO TAN, M-6

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Los Alamos National Laboratory
Los Alamos, New Mexico 87545

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PHASE DETONATED SHOCK TUBE (PFST)*

W. D. Zerwekh, S. P. Marsh and T.-H. Tan

Group M-6, Los Alamos National Laboratory, Los Alamos, NM 87545, USA

I. INTRODUCTION

The simple, cylindrically imploding and axially driven fast shock tube (FST) has been a basic component in our high velocity penetrator (HVP) program. It is a powerful device that is capable of delivering a directed and very high pressure output that we have successfully employed to drive hypervelocity projectiles. The FST is configured from a hollow, high explosive (HE) cylinder, a low density Styrofoam core, and a one point initiator at one end. A Mach stem formed in the core as the forward propagating, HE detonation wave intersects the reflected radial wave. By proper arrangement of HE length and diameter, a steady state Mach stem is readily achieved at the output end. Pre-detonation of HE is prevented by underdriving the Mach stem as it is being developed and this is most easily done by varying the foam density. The strength of the Mach stem is dependent on the effective energy transfer from the HF and this can be scaled geometrically. We have found this simple FST to be a powerful pressure multiplier. Typically, up to 1 Mbar output pressure can be obtained from this device. Further increase in the output pressure can be achieved by increasing the HF detonation velocity.

Over the last few years, the FST has been fine tuned to drive a thin plate to very high velocity under an impulse per unit area of about $1 \text{ Mbar} \mu\text{sec/cm}^2$. Typically a 1.0 mm thick, 1.0 to 1.5 cm thick, has been accelerated intact to $0.8 \text{ cm}/\mu\text{sec}$ under a loading pressure rate of several Mbar μsec . By making the plate on the side slightly convex at the loading side we have successfully accelerated it to about $1.0 \text{ cm}/\mu\text{sec}$. By placing a 1 mm layer of aluminum buffer, a thin titanium foil or titanium plate with an equivalent compressive mass, a plate has been accelerated to above

$1.1 \text{ cm}/\mu\text{sec}$. We have found the incorporation of a barrel at the end of the FST to be important. The confinement of the propellant gas by the barrel tends to accelerate the projectile to higher velocity. Furthermore, the standoff in the barrel between the plate and the FST allows the expanding gas to load more gently on the plate and thus reduces the loading pressure rate. A hookless acceleration is highly desirable to prevent the plate from being broken up prematurely. However, presence of a large standoff volume tends to introduce wall effects and generate serious perturbation from the not well understood, high pressure gas flow dynamics. We try to mitigate this difficulty by keeping the standoff distance as short as can be tolerated by experimental tests. In general we have found good agreement between the 2D numerical simulation and measurements. Even scaling tests appear to be satisfactory. A factor of three increase in the geometric dimensions of the FST, barrel, standoff, and plate yields similar results in both the calculations and experiments.

The desire to accelerate the plate above $1.0 \text{ cm}/\mu\text{sec}$ provides the impetus to develop a more advanced fast shock tube that will deliver a much higher output pressure. We decided to investigate a relatively simple air-burster phase detonation system (PEFD) with fifty percent higher phase detonation velocity and a modest 0.2 Mbar output. Code calculations show that the PEFD acceleration of a plate to about $1.2 \text{ cm}/\mu\text{sec}$ can be achieved. The performance of the ePEFD has been evaluated and the details are discussed below.

II. DESCRIPTION OF PHASE DETONATED FST

The phase detonated FST is shown in Figure 1a. The fast detonation velocity component of a Phase detonation HE burster detonates on the left of the bur-

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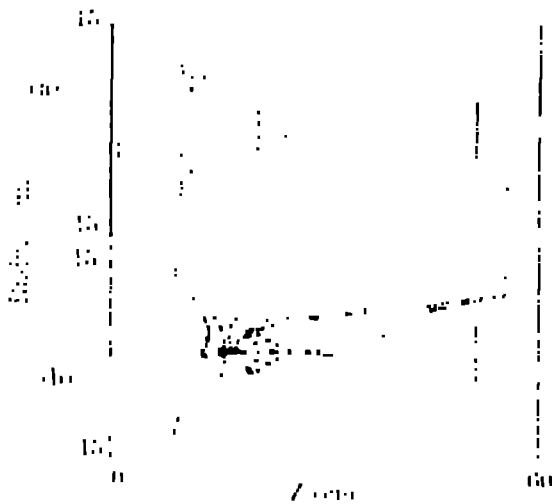


FIGURE 2
Phase detonated ESI with a bar, showing (a) collapse and (b) Mach disk formation process.

and the detonate the cylinder of Composition B. It is the outside surface of the ESI. As the detonation proceeds to the right, a cylindrical shell of 304 stainless steel SSU is propelled radially inward. It impinges on the conical surface of a 6061 aluminum phanor lens, the angle of which determines the phase velocity of the system (Figure 2). Thus, the transition of a shock to a cylindrical shell of PBX 9501 explosive and detonates it at an axial velocity determined by the phanor

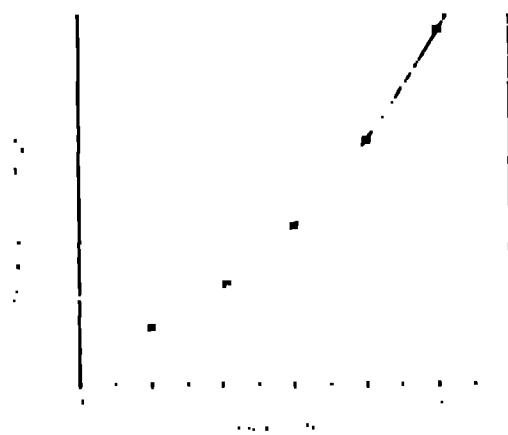


FIGURE 3
Calculated phase velocity as a function of cone angle for a PFI.

nor lens. An axial cylinder of Styrofoam is then shock compressed by this travelling detonation front resulting in a Mach disk travelling at the same phase velocity. It ultimately breaks out of the system at the right-hand face and causes the acceleration of a 304SS plate placed against that face or suspended in a 304SS barrel that rests against that face.

The purpose of the steel plate against the base of the PFI's lens, the polyethylene wedge against its periphery, and the polyethylene disk at the left face of the phanor enclosure to prevent the predetonation of the PBX 9501 explosive before the shock from the collapsing 304SS cylinder detonates it. The plug of polyethylene on the left face of the Styrofoam results in a delay of 100 distance for the full diameter formation of the Mach disk.

III. PHASE-DETONATION, MACH DISK PERFORMANCE

The hydrocode, Mea2D, has been used to model the overall performance of the phase detonated ESI. It is a 2D, Eulerian second order code that can handle multi-material problem and can treat the programmed detonation of high explosive. The results shown here have been obtained using a Cray YMP computer.

The modeling of a phase detonated ESI with a bar is shown in Figure 2. The inward collapse of the 304 stainless steel on to the 6061 Al phanor lens, shock in the lens, phase detonation of the PBX 9501 explosive and Mach disk formation. The angle of the phanor lens in this problem is 45 degree and the resulting phase velocity of the PBX 9501 detonation front is 1.40 cm/μs, an increase of 50% over the normal detonation velocity of 0.9 cm/μs. The shock velocity of the Mach disk in the 304SS barrel is the same value of 1.40 cm/μs.

The one-dimensional experiment has been performed to measure the phase velocity in the system. The measurements to be accomplished are: (a) the Mach disk formation, (b) axial cable zero embedded at equal distance interval around an aluminum rod (c) pressure measurement probe that replaced the face of 200 holes (d) shock absorber to remove the wave reflection. The results are all computer simulated reaction

conventional process machined into the aluminum rod and their lead protruding outer conductors carefully re-spooled, exposing the central coaxial conductor (Figure 3). As the phase detonation proceeded along the PBX 9501 explosive, the cables were shunted and through an interferometry circuit the position vs. time data of the two cables was determined. The velocity was then extracted from the x-t data.

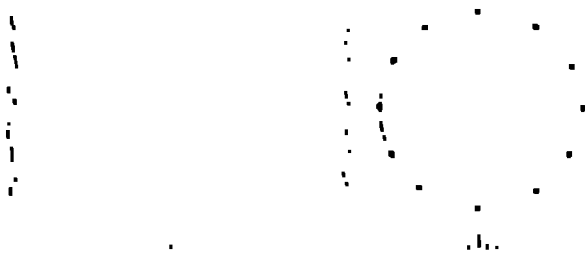


FIGURE 3

Side-on and end-on views of microwave interferometer probe cables to determine phase velocity.

A plot of the detonation velocity average of the channels is shown in Figure 4. The phase velocity is approximately 1.49 cm/μs, which is acceptable, close to the design value of 1.50 cm/μs. The shot was visually inspected.

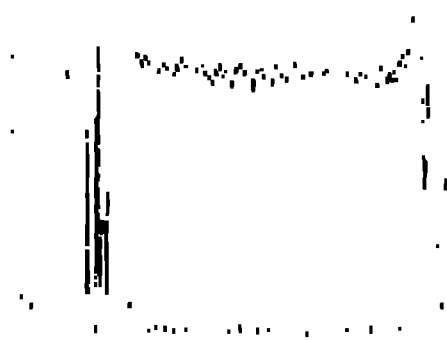


FIGURE 4

Average detonation velocity of the interferometer probe cables.

IV. PLATE ACCELERATION PERFORMANCE

Because of the poor results obtained when trying to accelerate a plate down a barrel it was decided to attempt (with a vertical axis system) to accelerate a plate initially in contact with the shock tube face and having convex plate curvature toward the shock tube face as shown in Figure 5. The plate is 0.15 cm thick, 3018S and for the two systems studied the radii of curvature of the plate faces against the shock tube were 1130 cm and 672 cm. Hydrocalculation in these two systems are shown in overlays on Figure 5. The effect of decreasing the radius of curvature is evident in the greater consolidation, thickness, and increased forward bowing. The plate velocities are 1.26 and 1.20 cm/μs for the low and the high radius of curvature plates, respectively.

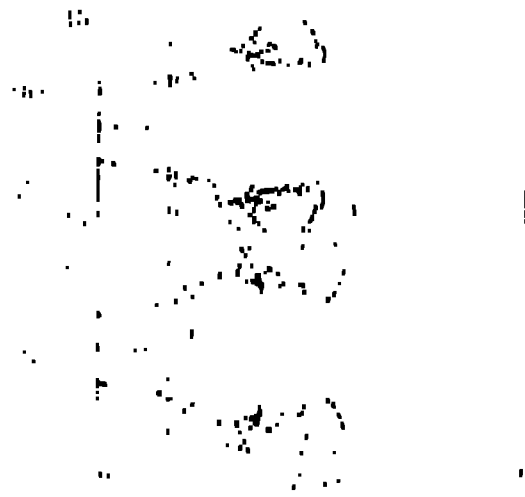


FIGURE 5

Overlays of calculated plate contour at 1 μs in two systems for phase detonated 1.49 cm/μs barrel. The radii of curvature of the plate are 1130 cm and 672 cm.

Experiment was performed on shock tube with the convexity of curvature machined on the plate surface. A hydrophone was used to help in comparison of the systems and PBX 9501. The velocity data and pressure data of curvature of 1130 cm and 672 cm are shown in Figure 6.

The experimental setup for these shots is shown in Figure 6. The assemblies accelerated the plates vertically downward. The displacements of the plates at the 100- μ s ray times (9.5 and 36.1 cm) were greater than the largest displacements shown in the hydrocode simulation (6.3 cm in both calculations) in order to protect the x-ray head and the film cassette from blast and impact.

The radiographic results for these shots are given in Figure 7. Only one exposure was obtained for shot 57-476 because only one channel of the film's x-ray unit was triggered. But a plate velocity was obtained for this shot by time information on initial motion time obtained from hydrocode simulations and the earlier shot fired. The velocity of the plate in shot 57-476 is 1.18 cm μ s⁻¹ and the velocity of the negative part of the tip in shot 57-476 is 1.16 cm μ s⁻¹.

The difference in the calculated and experimental plate contours for the plate with 1137 cm radius of curvature (figure 7a) and 7a(i) caused by early fracture of the plate at its periphery (not modeled by the hydrocode reducing the driving pressure in that region and causing it to lag behind the plate center, which can be seen in Figure 6). Although the same peripheral pressure drop exists for the plate with the smaller radius of curvature, the convergence is much greater that it folds forward and collapses on itself as the side displacement is greater than the displacement needed to meet it. The plate in both of these shots may be related to the residual internal energy after the initial

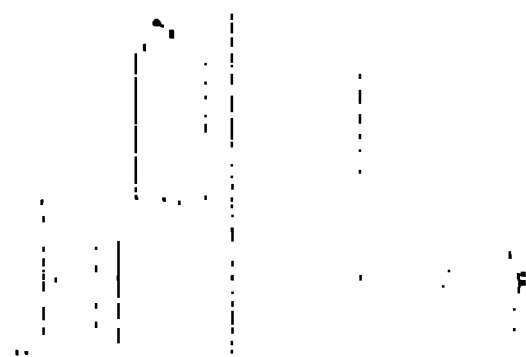


FIGURE 6
Experimental setup for shots 57-476 and 57-475.

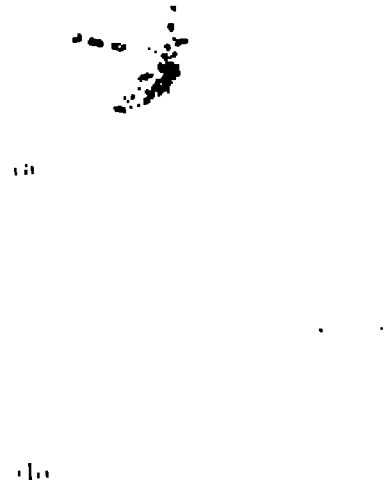


FIGURE 7
Radiograph of thin steel plate, moving to the right in shot 57-476 and (b) shot 57-475. Over-constriction of plate in (b) can be seen to have occurred.

burst and detonation are completed and the plate is back at zero pressure.

V. CONCLUSIONS

Detonation of PBX 9402 at a plate velocity of 1.10 cm μ s⁻¹ has been observed, and the detonation velocity was uniform at 19 positions around the inner cylindrical interface. A Mach disk in an axial symmetry was formed within the explosive and traveled at the same plate velocity. Pressure fluctuations caused by the foam-explosive interface were eliminated by firing the shot vertically and placing a low velocity area in the interface gap. A radius of curvature of 1137 cm on the driving face of a 7% plate in contact with the foam III face of a back tube produced plate convergence during acceleration and allowed the plate to reach a velocity of 1.1 cm μ s⁻¹. The only one edge of the plate indicates it may have melted. Arrangement with a barrel of appropriate length will allow the use of detonation to initiate a plate for shock and impact.